

## Radiation Balance for Earth

### Recap:

- Solar radiation at TOA is ~ blackbody at 5780 K
- Has small solid angle  $\rightarrow$  plane wave
- Integrated energy on a normal to the surface is  $\sim 1370 \text{ W/m}^2$
- varies with latitude and season

We can calculate effective T of Earth by considering radiative equilibrium between solar and terrestrial radiation.

### Energy absorbed by Earth

= energy intercepted  $\times$  (1 - average reflection coefficient)

$$= \pi R^2 F (1 - A)$$

where

$R$  = Earth's radius

$F$  = solar flux density at Earth

$A$  = planetary albedo (= ratio of reflected flux density to incoming solar flux density)

Energy radiated by Earth

= emissivity  $\times$  surface area  $\times$  blackbody flux density

$$= \epsilon 4\pi R^2 \sigma T_e^4$$

$\hookrightarrow \approx 1$  for infrared

$$\therefore 4\pi R^2 \sigma T_e^4 = \pi R^2 F(1-A)$$

$$T_e = \left[ \frac{F(1-A)}{4\sigma} \right]^{1/4}$$

defines the effective radiating temperature of Earth (or any planet)

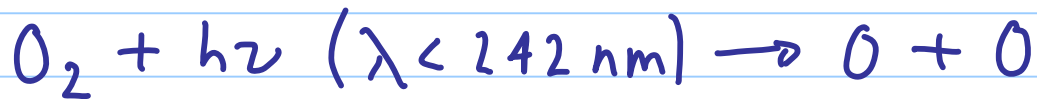
For Earth,  $T_e = 255 \text{ K}$

The atmosphere has a warming effect that makes the Earth's surface temperature about 33 K warmer than  $T_e$ .

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## Oxygen and Ozone in the Atmosphere

- ① atomic oxygen created by photolysis of molecular oxygen



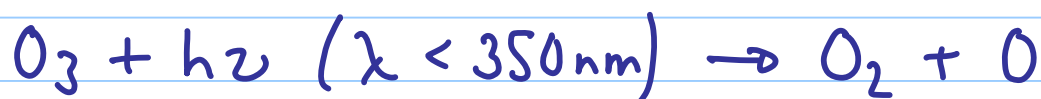
photolysis rate  $J_1$

- ② production of ozone



rate coefficient  $k_1$

- ③ photolysis of ozone



photolysis rate  $J_2$

- ④ collision of ozone with atomic oxygen



rate coefficient  $k_2$

- ⑤ self reaction of atomic oxygen



} very slow,  
can ignore  
< 60 Km

These reactions occur simultaneously  
 → Chapman Cycle

Reactions (2) and (3) are faster than (1) and (4), so for equilibrium of O and O<sub>3</sub> (odd oxygen):

$$k_1 [O][O_2][M] = J_2 [O_3]$$

$$k_2 [O][O_3] = J_1 [O_2] \text{ for overall balance.}$$

Combine these:

$$[O] = J_1 [O_2] / k_2 [O_3]$$

$$\begin{aligned} [O_3] &= \frac{k_1}{J_2} [O][O_2][M] \\ &= \frac{k_1}{J_2} \frac{J_1 [O_2]}{k_2 [O_3]} [O_2][M] \end{aligned}$$

$$\therefore [O_3] = [O_2] \sqrt{\frac{J_1 k_1 [M]}{J_2 k_2}}$$

So [O<sub>3</sub>] depends on J<sub>1</sub>/J<sub>2</sub>, which increases with height.

$$J = \int_{\lambda} \sigma_{\lambda} F_{\lambda}(\infty) \tau_{\lambda}(z, \infty) d\lambda$$

→ dissociating quanta per molecule that are absorbed

$[O_3]$  also depends on  $[M]$  and  $[O_2]$  which decrease with height.

∴ There is maximum in  $[O_3]$  at some height

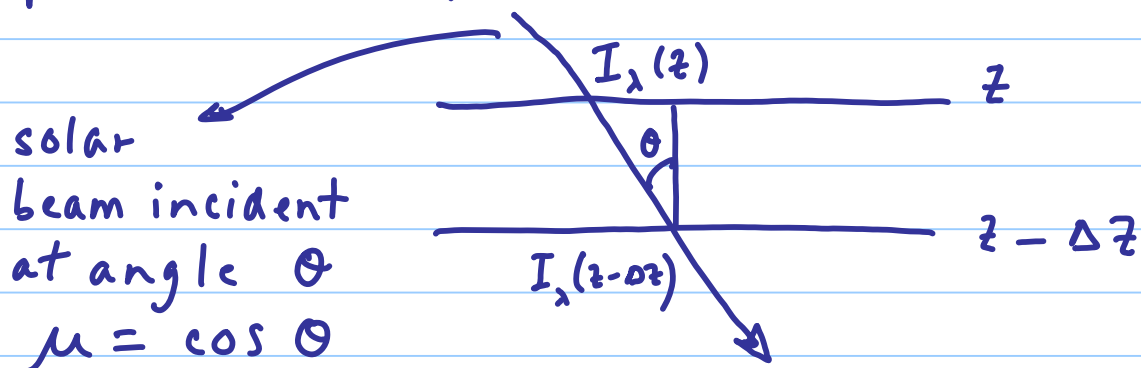
→ ozone layer

There will also be a maximum in heating of the atmosphere at this height

→ stratosphere

## Solar Heating of the Atmosphere

Consider an atmospheric layer in a plane-parallel atmosphere.



$$\begin{aligned}
 I_{\lambda}(z - \Delta z) &= I_{\lambda}(z) \tau_{\lambda}(\Delta z / \mu) \\
 &= I_{\lambda}(z) \exp\left(-k_{\lambda} \rho \frac{\Delta z}{\mu}\right)
 \end{aligned}$$

using

Schwarchild's Equation for a homogeneous layer with  $J=0$ .

Energy lost (per unit area)

$$\begin{aligned}
 &= [I_{\lambda}(z) - I_{\lambda}(z - \Delta z)] \Delta \Omega_s \\
 &= [I_{\lambda}(z) - I_{\lambda}(z) \tau_{\lambda}(\Delta z / \mu)] \Delta \Omega_s \\
 &= I_{\lambda}(z) A(\Delta z / \mu) \Delta \Omega_s \quad \text{where } A_{\lambda} = 1 - \tau_{\lambda} \\
 &= \frac{\Delta E_{\lambda}}{\Delta t}
 \end{aligned}$$

This energy loss from the beam is also the net energy gain of the atmosphere, so

$$\frac{\Delta E_{\lambda}}{\Delta t} = \rho C_p \Delta V \frac{dT}{dt}$$

where  $\rho$  = density

$C_p$  = specific heat at constant pressure

$\Delta V$  = volume of cylinder where heat is absorbed

→ J/kg/K

$\Delta V = \Delta z / \mu$  assuming a cylinder of unit cross-section aligned along the beam

$$\begin{aligned} \frac{dT}{dt} &= \frac{1}{\rho C_p \Delta z / \mu} \frac{\Delta E}{\Delta t} \\ &= \frac{\mu}{\rho C_p \Delta z} I_\lambda(z) A_\lambda (\Delta z / \mu) \Delta \Omega_s \end{aligned}$$

From Hydrostatic Egn  $\Delta p = -\rho g \Delta z$

$$\frac{dT}{dt} = \frac{-g}{C_p} \frac{\mu I_\lambda(p) A_\lambda (\Delta p / \mu) \Delta \Omega_s}{\Delta p}$$

$$\frac{dT}{dt} = -\Gamma \frac{dF_\lambda(\mu, p)}{dp} = \underline{\text{heating rate}}$$

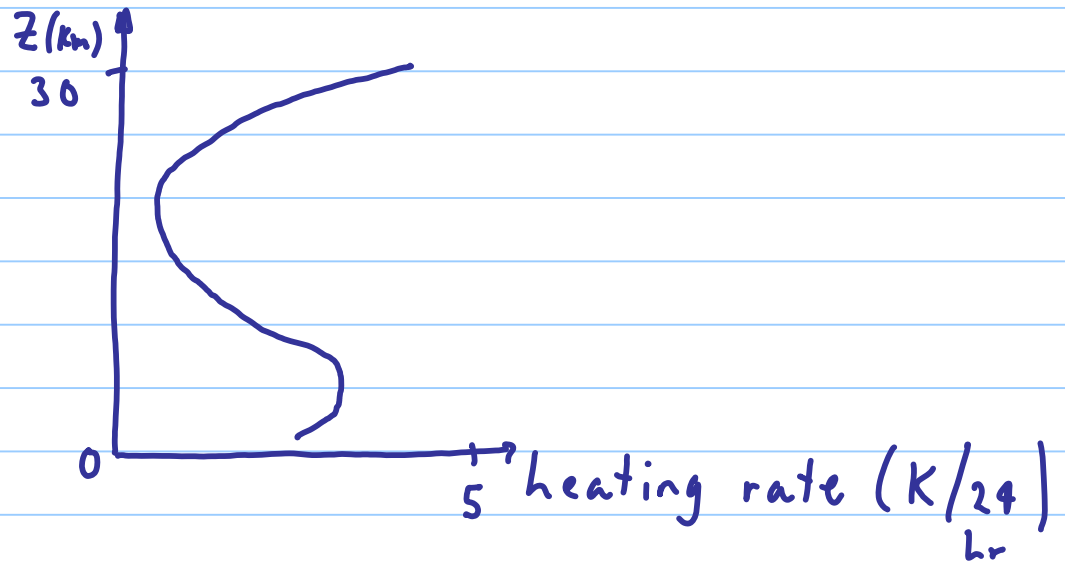
where

$\Gamma = g / C_p =$  adiabatic lapse rate

$F_\lambda = \mu I_\lambda(p) A_\lambda (\Delta p / \mu) \Delta \Omega_s =$  radiant flux density incident on unit area

Since  $dF_\lambda / dp$  is always  $< 0$ , so  $\frac{dT}{dt}$  is  $> 0$ .

This equation allows calculation of the heating rates in the atmosphere due to solar radiation.



Plotting the heating rate for solar radiation in the lower atmosphere in  $^{\circ}\text{C}/\text{day}$  shows:

- peak at the ground due to  $\text{H}_2\text{O}$  and  $\text{CO}_2$  absorption
- then a decrease to a minimum
- above the tropopause,  $\text{O}_3$  absorption becomes important, causing a large increase in heating rate

The above calculation of  $dT/dt$  is for monochromatic radiation.



Total heating rate due to solar radiation is

$$\left(\frac{dT}{dt}\right)_{\text{total}} = \sum_{\lambda} \left(\frac{dT}{dt}\right)_{\lambda}$$

$$= -\Gamma \sum_{\lambda_i} \left. \frac{dF_{\lambda_i}}{dp} \right|_p \Delta\lambda_i$$

sum  
over all  
spectral  
intervals  
'i'