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equatorial jets in the ocean (motivation)

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- applications

Observations of the equatorial Atlantic:

5°S











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- symmetry about the equator in both zonal and meridional velocity (Bourles et al., 2003).
- Dispersion diagram for meridionally confined equatorial waves:
 - Short wavelength Rossby and mixed Rossby-gravity waves have low-frequency and eastward group velocity.
 - Variability near western boundary could excite such waves which would then propagate into the interior.



Jet formation through wave instability

In the short wavelength limit, the instability of barotropic Rossby waves (Lorenz, 1972; Gill, 1974; Lee and Smith, 2003) leads to the formation of zonal jets (Manfroi and Young, 1999).

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- Equatorial Rossby and mixed Rossby-gravity (MRG) waves have dispersion relations and velocity fields similar to barotropic Rossby waves and are thus subject to the same barotropic instability (even though they are inherently baroclinic).

- Numerical simulations of MRG waves excited from the western boundary show that they can lead to zonal jets:
 - Baroclinic equatorial deep jets
 (d'Orgeville et al., 2007; Hua et al., 2008)
 - Extra-equatorial barotropic jets (Ménesguen et al., submitted)



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- giving:
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- Parameters varied were the amplitude of meridional velocity V_0^* and the zonal wavelength λ_x , or in nondimensional terms:

• Froude number
$$V_0 \equiv \frac{V_0^*}{c}$$
, zonal wavenumber $k = 2\pi \frac{L_D}{\lambda_x}$

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 \Rightarrow zonal shear in meridional velocity dominates the β effect (rotation) and ω (wave propagation).

- Simulation with $V_0 = 0.11$ (0.36 cm amplitude) and k = -6.3 (350 km wavelength)
 - at $0.1^{\circ} \times 0.1^{\circ}$ horizontal resolution ($1^{\circ} \approx 100$ km)
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- A necessary condition for instability is that the flow V(x) have an inflection point (Rayleigh), and furthermore, that it be a maximum of absolute vorticity (Fjortoft).



• "Rossby" waves propagate in the positive y direction relative to background flow where $V_{xx} > 0$, and in the negative y direction where $V_{xx} < 0$.

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For small enough values of l, waves can phase lock and grow exponentially.

Scale length and time like $(x, y) = k^{-1}(\xi, \eta)$ and $t = (kV_0)^{-1}\tau$. Basic state wave is:

$$V \sim V_0 \exp\left(\frac{-\eta^2}{2k^2}\right) \cos(z) \cos(\xi),$$

$$U \sim k^{-1} V_0 \left(\frac{\eta}{k}\right) \exp\left(\frac{-\eta^2}{2k^2}\right) \cos(z) \sin(\xi),$$

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$$\tilde{\zeta'}_{\tau} + \underbrace{V\tilde{\zeta'}_{\eta}}_{\text{adv. of }\zeta' \text{ by }V} + \underbrace{k^{-1}Z_{\xi}u'}_{\text{adv. of }Z \text{ by }u'} = \underbrace{\frac{V_0}{k^2}(UZ_{\xi} + VZ_{\eta})}_{\text{self-interaction of wave}}$$

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Next order in k^{-1} : time dep. of MRG wave, advection by U, β effect.

Look for horizontally non-divergent perturbations $u' = -\psi'_{\eta}$, $v' = \psi'_{\xi}$. Then

$$\nabla^2 \psi'_{\tau} + \cos(z) \exp\left(\frac{-\eta^2}{2k^2}\right) \cos(\xi) \left(\nabla^2 \psi'_{\eta} + \psi'_{\eta}\right) = \frac{V_0}{k} S[\cos(2\xi), \sin(2\xi), y, z]$$

Linear partial differential equation with coefficients periodic in ξ .
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Seek solution with periodicity of coefficients (Floquet theorem with exponent $n_0 \equiv 0$).

$$\psi'(\xi,\eta,z,t) = \Re\left\{e^{in_0\xi}\sum_{n=-\infty}^{\infty}\hat{\psi}_n(\eta,t;z)e^{in\xi}\right\}$$

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- Coefficient proportional to $\cos(\xi)$ couples adjacent terms in the summation.
- $\psi_0(\eta, t; z)$ corresponds to the zonally symmetric zonal velocity (the zonal jet).
- A good approximation is obtained with a three term truncation:

$$\psi' \approx \Re \left\{ \hat{\psi}_{-1}(\eta, z, t) e^{-i\xi} + \hat{\psi}_0(\eta, z, t) + \hat{\psi}_1(\eta, z, t) e^{i\xi} \right\}$$

Substituting the truncated series into the PDE, setting the coefficients of $e^{-i\xi}$, e^0 , and $e^{i\xi}$ equal to zero, and discretizing the differential operators in η yields:

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Right hand side S acts like an initial seed for the perturbation.

Growth rate is proportional to $\cos(z)$.



Most unstable mode for k = -6.3:



even and odd modes come in pairs, with odd modes having slightly higher growth rate.



- Jet spacing near equator matches barotropic Rossby wave case (Gill, 1974).
- Jet spacing wider than wavelength of MRG wave l < 1.
- Fastest growing jet has width 1.6 times zonal scale of basic state wave.





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- Equatorial jet always westward. Probably due to mixing of planetary angular momentum by zonally short modes (strongest at equator) biasing simulations towards westward flow at equator.
- Extra-equatorial jet positioning poleward of that in most unstable linear mode: barotropic instability?



barotropic instability possible when $U_{yy} = \beta$ (does not lead to meridional jets due to Coriolis effect).

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 - It forces the meridional phase of the jets to be such that a jet is centred on the equator.
 - It restricts the maximum amplitude of the equatorial westward jet.
 - It leads to latitudinal homgenization of angular momentum and potential vorticity in the vicinity of the equator.
 - It leads to vertical mixing of density and tracers in the equatorial track.







- Balanced centrifugal $\frac{M^2}{r^3}$ and pressure gradient $-\frac{\partial P}{\partial r}$ forces.
- Consider a ring of fluid at r_1 displaced to r_2 conserving its angular momentum without disturbing the pressure field.



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Consider a ring of fluid at r_1 displaced to r_2 conserving its angular momentum without disturbing the pressure field. At its new position, it feels a net radial force

$$\frac{M_1^2}{r_2^3} - \left. \frac{\partial P}{\partial r} \right|_{r_2} = \frac{M_1^2 - M_2^2}{r_2^3} \approx \left. -\frac{1}{r_2^3} \left. \frac{\partial \left(M^2 \right)}{\partial r} \right|_{r_2} \delta r$$



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 \Rightarrow An axially symmetric flow is inertially stable if $\frac{\partial (M^2)}{\partial r} > 0$ (Rayleigh, 1917).



A famous example of inertial instability is the instability of laminar flow in the Taylor-Couette experiment, when outer cylinder rotates too fast relative to inner cylinder (Taylor (1923)).



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- First bifurcation leads to Taylor vortices in the radial-vertical plane.
- For inertial stability on the earth, maximum angular momentum must be at the equator.
- Combined with effects of stratification, the condition for stability is that Potential Vorticity have the sign of latitude ("symmetric stability").

For stability, maximum angular momentum $M \approx U - \frac{1}{2}\beta y^2$ must be at the equator:

Equatorial shear inertial instability



"Curvature" inertial instability



(Dunkerton, 1981)

 $U = \lambda y$, unstable for all $\lambda \neq 0$

Adjustment mixes in latitude and depth.

$$\begin{split} U \;&=\; (-U_{00} \;+\; \frac{1}{2} b y^2) e^{-\frac{1}{2} y^2},\\ \text{stable for } b \;\!>\;\!\beta \end{split}$$

un-

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inertial adjustment takes place over wider latitude interval than covered by initial instability (Hua et al., 1997, Griffiths, 2003).

Results in wide area of zero PV.

Inertial instability in simulations

Evidence of inertial instability in the simulations:

Absolute vorticity vs. latitude



Also, energy increase in zonally symmetric components of v and w fields.

Jet amplitude and anomalous PV

Baroclinic equatorial deep jets

If the instability is barotropic, why are the jets seen in the observations and in basin simulations (d'Orgeville et al., 2007) alternating in sign with depth?



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Hua et al., 2008:

- Since the destabilization is taking place mostly in the western edge of the basin, what is observed in the remainder of the basin is the projection of the jets pattern onto low-frequency, long zonal wavelength waves with eastward group velocities.
- The only such wave is the equatorial Kelvin wave, which has

$$u(y,z) \propto \cos(mz) \exp\left[-\left(\frac{\beta m}{2NH}\right)y^2\right]$$

• The vertical mode m of the observed jets corresponds to the vertical mode of the Kelvin wave with meridional width $\sqrt{NH/\beta m}$ comparable to the width of the most unstable barotropic jet mode, namely $m \propto k^2 \propto T^2$.



Layering, zonal jets, and inertial instability

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(Ménesguen et al., submitted)

Regions of vertical mixing are concentrated in westward jets, and are strongly correlated with the criterion for marginal inertial instability.

Summary

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- Baroclinic Equatorial Deep Jets predominate in basin simulations because they correspond to Kelvin waves, which have eastward group velocities. (Hua et al., 2008)
- Inertial instability leads to vertical mixing of density and tracers in equatorial track. (Ménesguen et al., submitted)

Continuing Work

- Very short wavelength waves lead to local super-rotation. Application to planetary atmospheres?
- Baroclinic effects of depth-dependent triggering of inertial instability.
 - The effect of the non-traditional
 Coriolis force? Simulations show strong vertical symmetry breaking.
 Evidence of global conservation of angular momentum for wave momentum flux?

