

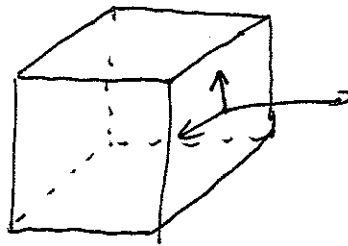
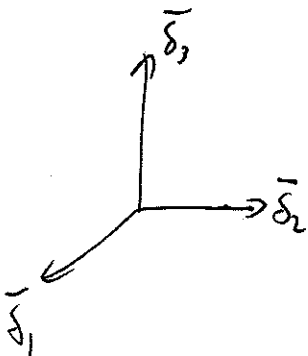
Tensor Review

Scalar \rightarrow specified by magnitude (0th rank tensor)

vector \rightarrow magnitude and direction (1st rank tensor)

\Rightarrow tensor specifies magnitude and a number of directions. ($n \rightarrow n$ th rank tensor)

Stress tensor = 2nd rank tensor



$$\tau = \begin{pmatrix} \tau_{11} & \tau_{12} & \tau_{13} \\ \tau_{21} & \tau_{22} & \tau_{23} \\ \tau_{31} & \tau_{32} & \tau_{33} \end{pmatrix}$$

Vector dot product

$$\vec{u} = u_1 \vec{\delta}_1 + u_2 \vec{\delta}_2 + u_3 \vec{\delta}_3$$

$$\vec{v} = v_1 \vec{\delta}_1 + v_2 \vec{\delta}_2 + v_3 \vec{\delta}_3$$

$$\vec{u} \cdot \vec{v} = \sum_i u_i \vec{\delta}_i \cdot \sum_j v_j \vec{\delta}_j = \sum_i \sum_j (u_i v_j) (\vec{\delta}_i \cdot \vec{\delta}_j)$$

$$\text{Kronecker Delta} \equiv (\vec{\delta}_i \cdot \vec{\delta}_j) = \delta_{ij}$$

$$\Rightarrow \boxed{\bar{u} \cdot \bar{v} = u_i v_i}$$

Einstein summation convention: repeated indices imply summation.

Cross product :

$$\bar{u} \times \bar{v} = (u_2 v_3 - u_3 v_2) \bar{e}_1 + (u_3 v_1 - u_1 v_3) \bar{e}_2 + (u_1 v_2 - u_2 v_1) \bar{e}_3$$

introduce permutations symbol

$$\epsilon_{ijk} = \begin{cases} +1 & ijk \text{ cyclic order (eg } 123, 231, 312) \\ 0 & \\ -1 & ijk \text{ anticyclic order (eg } 321, 213, 132) \end{cases}$$

$$\Rightarrow \boxed{\bar{u} \times \bar{v} = u_i v_j \epsilon_{ijk} \bar{e}_k}$$

$$\begin{aligned} \epsilon_{ijk} &= \epsilon_{jki} = \epsilon_{kij} \\ \epsilon_{ijk} &= -\epsilon_{ikj} \end{aligned}$$

$$\boxed{\bar{e}_i \times \bar{e}_j = \epsilon_{ijk} \bar{e}_k}$$

Second rank tensor

$$\bar{T} = T_{11} \bar{e}_1 \bar{e}_1 + T_{12} \bar{e}_1 \bar{e}_2 + T_{13} \bar{e}_1 \bar{e}_3 + \dots$$

$$= T_{ij} \bar{e}_i \bar{e}_j$$

$\bar{e}_i \bar{e}_j = \text{"Dyadic tensor"}$

$$\boxed{\begin{aligned} \bar{u} \bar{v} &= u_i \bar{e}_i v_j \bar{e}_j \\ &= u_i v_j \bar{e}_i \bar{e}_j \end{aligned}}$$

(3)

unit tensor $\bar{\bar{I}} = \delta_{ij} \bar{\delta}_i \bar{\delta}_j$

Operations on tensors

1) $\bar{\delta}_i \bar{\delta}_j \cdot \bar{\delta}_k \bar{\delta}_l = \delta_{jk} \bar{\delta}_i \bar{\delta}_l \Rightarrow$ tensor

2) $\bar{\delta}_i \cdot \bar{\delta}_j \bar{\delta}_k = \delta_{ij} \bar{\delta}_k \Rightarrow$ vector

3) $\bar{\delta}_i \bar{\delta}_j : \bar{\delta}_k \bar{\delta}_l = \delta_{jk} \delta_{il} \Rightarrow$ scalar

(double dot product)

4) $\bar{\bar{I}} : \bar{\bar{A}} = A_{ii} \Rightarrow$ trace of A

$\bar{\bar{I}} : \bar{\bar{I}} = 3$

Calculus of vectors and Tensors

Divergence: 1) $\nabla \cdot \bar{u} = \frac{\partial}{\partial x_i} \bar{\delta}_i \cdot u_j \bar{\delta}_j = \frac{\partial u_j}{\partial x_i} \delta_{ij}$
 $= \frac{\partial u_i}{\partial x_i}$

2) $(\nabla \cdot \bar{\bar{T}})_i = \frac{\partial T_{ij}}{\partial x_j}$

3) curl

$$\nabla \times \vec{u} = \frac{\partial}{\partial x_i} \vec{e}_i \times u_j \vec{e}_j = \frac{\partial u_j}{\partial x_i} (\vec{e}_i \times \vec{e}_j)$$

$$= \frac{\partial u_j}{\partial x_i} \epsilon_{ijk} \vec{e}_k$$

If $\nabla \cdot \vec{u} = 0$ the field is called solenoidal

$\nabla \times \vec{u} = 0$ the field is called irrotational.

Symmetric and Antisymmetric tensors

1) A tensor is symmetric if $B_{ij} = B_{ji}$

2) antisymmetric if $B_{ij} = -B_{ji}$

3) A tensor can be represented as a sum of a symmetric and antisymmetric tensor

$$B_{ij} = \underbrace{\frac{1}{2}(B_{ij} + B_{ji})}_{\text{symmetric}} + \underbrace{\frac{1}{2}(B_{ij} - B_{ji})}_{\text{antisymmetric}}$$