

# Influence of the Stratospheric Potential Vorticity Distribution on the Brewer-Dobson Circulation

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Brewer-Dobson circulation is driven in part by breaking of planetary scale Rossby waves on the winter polar vortex.

Both wave propagation and wave breaking are controlled by PV distribution  
(two separate ideas - linear, nonlinear)

e.g.

wave propagation:

- . . . steep gradients act as a waveguide
- . . . different dispersion relation for edge waves (polar vortex edge) and plane waves (uniform PV gradient)

wave breaking:

- . . . steep PV gradients limit wave breaking

Know that winter stratospheric PV has step-like distribution (vortex edge, surf-zone, subtropical barrier)

. . . aircraft observations ([Tuck, 1989](#); [Vaugh et al., 1994](#))

. . . dynamics of vortex stripping well-understood ([Legras et al., 2001](#))

. . . also high resolution re-analysis; contour advection; etc

To resolve this PV distribution requires high model resolutions.

Low resolutions will smooth the step.

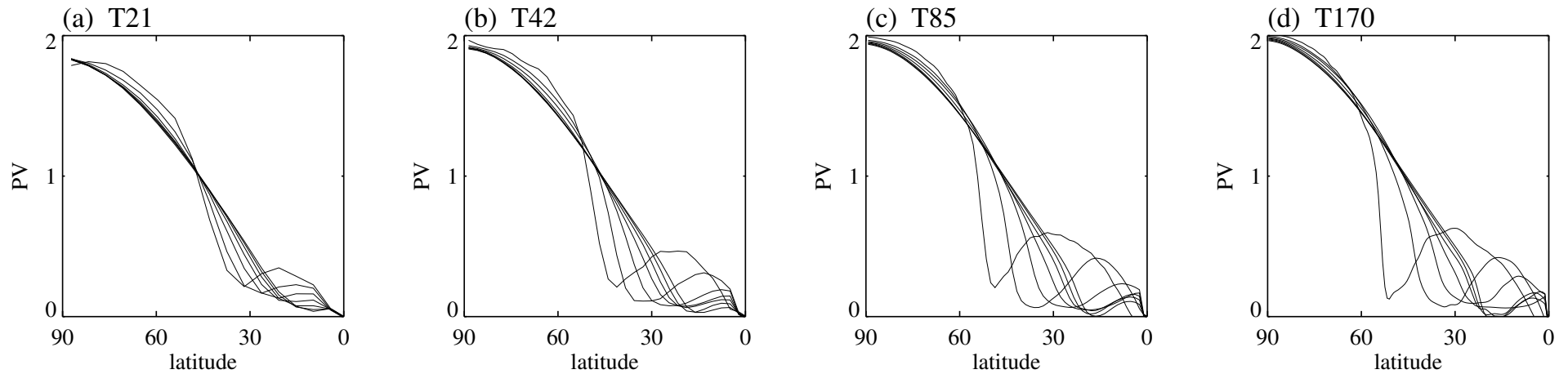
(e.g. T21: 16 grid points pole-equator; vortex edge width is 20 degrees)

. . . this will affect wave propagation/breaking

. . . diffusion + strain leads to anomalous vortex erosion and spurious transport.

**How important is this?**

Example: cross section of PV on  $\theta = 930$  K for T21 – T170



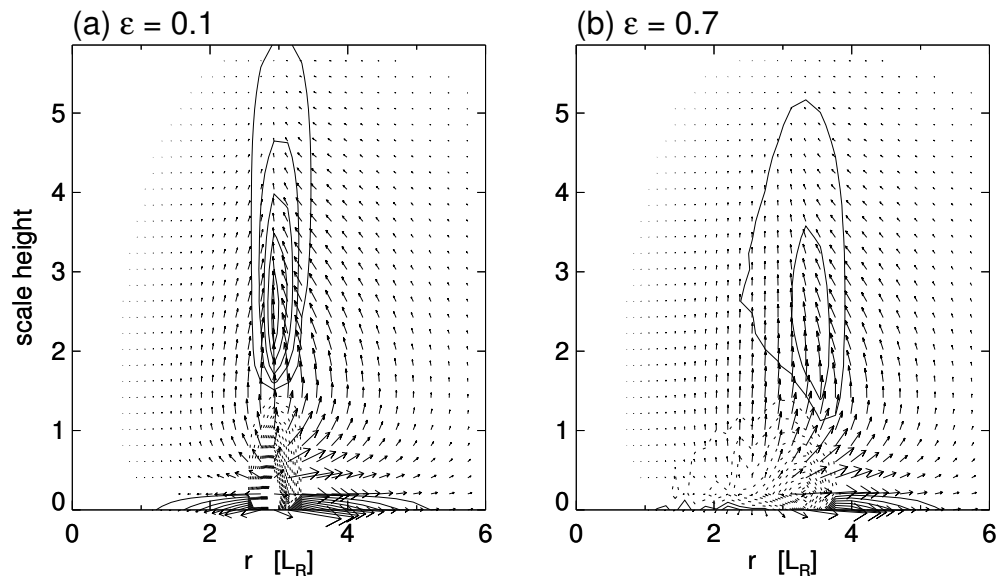
- . . . identical forcing and initial conditions (smooth vortex)
- . . . lines are  $t = 16, 18, 20, \dots, 28$  days
- . . . sudden warming around  $t = 28$  days

Scott et al. (2004):

quasigeostrophic model; vortex edge width  $\varepsilon$ ;

vortex parameters  $\kappa$ ,  $J$ ,  $r_0$  independent of  $\varepsilon$

modest increase in  $\nabla \cdot \mathbf{F}$  with decreasing  $\varepsilon$



But:

(i) difference in  $\nabla \cdot \mathbf{F}$  partly due to change in  $F^{(z)}$  at lower boundary  
(controlled by stratosphere)

(ii) did not consider barotropic mode

# INTERNAL STRATOSPHERIC VARIABILITY

Steady forcing → *coherent, regular* stratospheric variability

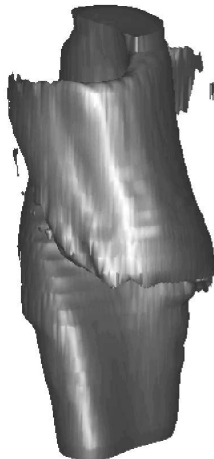
Holton & Mass; Yoden; Scott & Haynes; Christiansen

. . . shown to persist at regular internal variability at higher resolutions  
(Scott & Polvani, 2004, 2006)

(d)  $t = 690$



(e)  $t = 695$



(f)  $t = 700$



natural variability on timescales of  
around 60 days

repeated sudden warmings

competition between dynamical forcing  
(wave breaking) and thermal relaxation

. . . pseudo-spectral, stratosphere-only primitive equation model

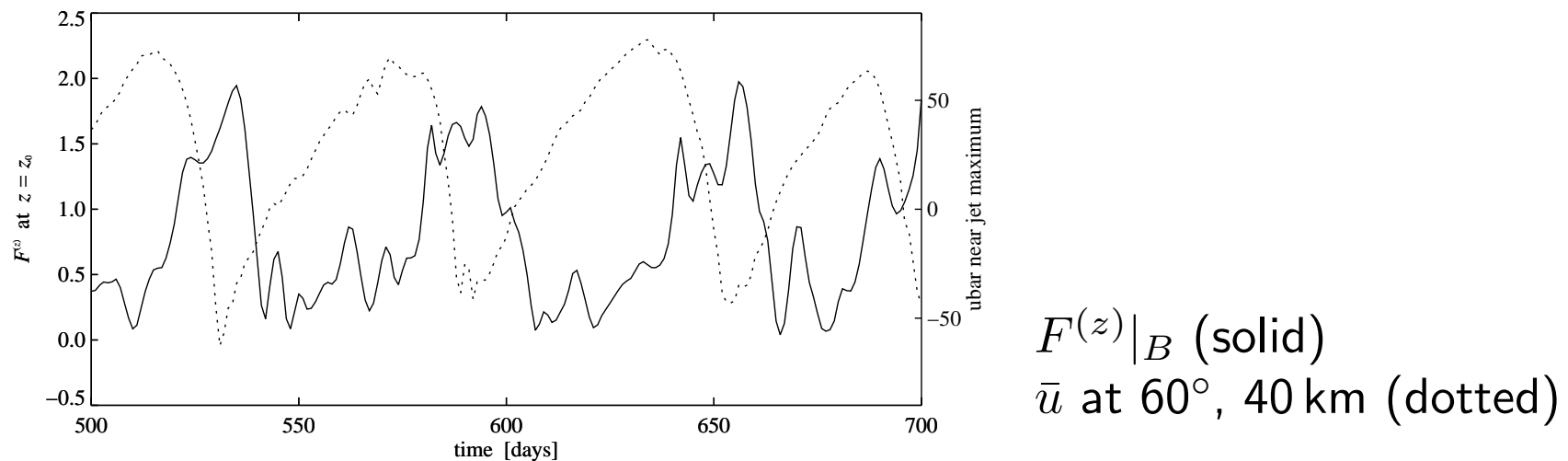
. . . 40 vertical levels

. . . thermal forcing to perpetual January

. . . lower boundary wave forcing – geopotential height perturbation

Here, wave forcing is time-independent (no tropospheric variability), yet *wave flux* through the lower boundary is strongly time-dependent: rapid increase as the polar vortex is decelerated

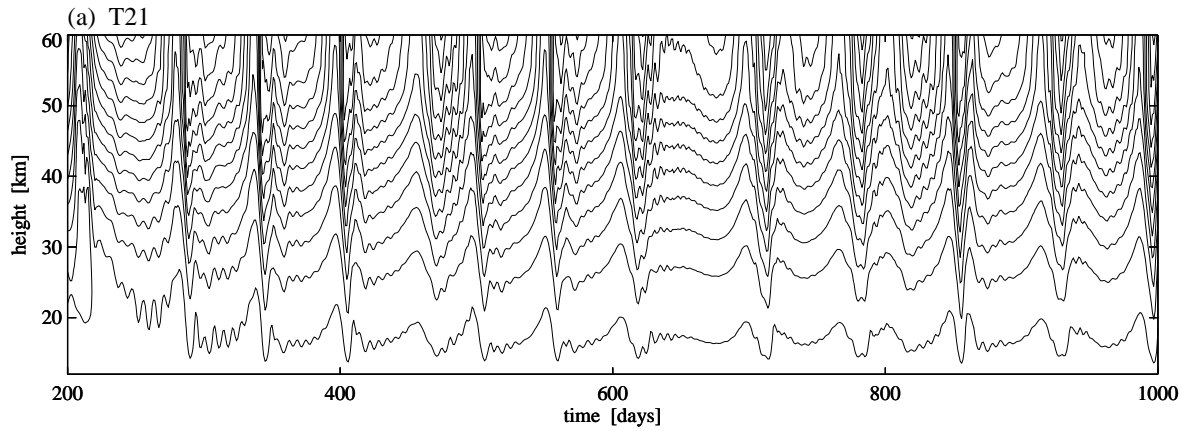
Sudden variations in  $F^{(z)}$  are determined by the state of the stratosphere (Scott & Polvani, 2004, 2006)



*A sudden increase in wave flux at the tropopause is not necessarily caused by an increase in tropospheric wave forcing.*

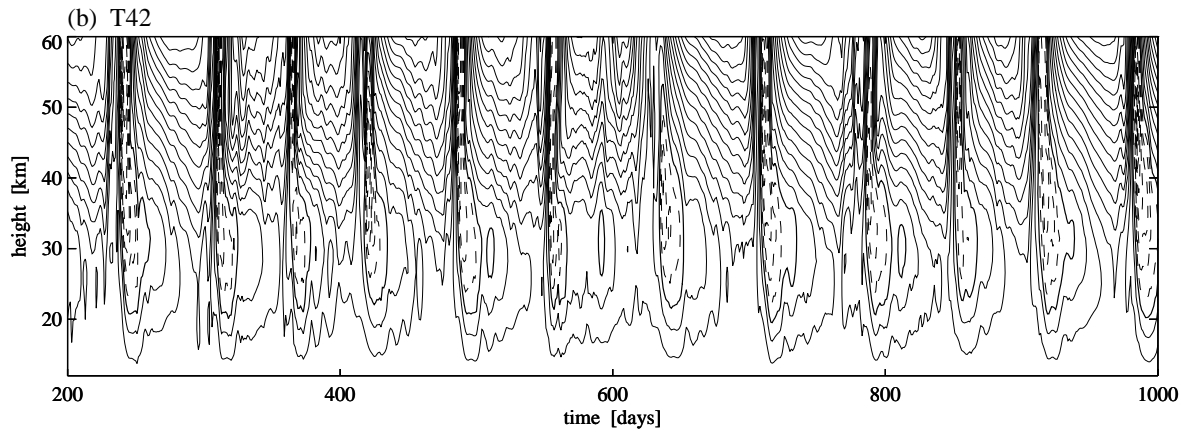
*The stratosphere alone is able to modulate the wave flux.*

# sensitivity to horizontal resolution: T21, T42, T85

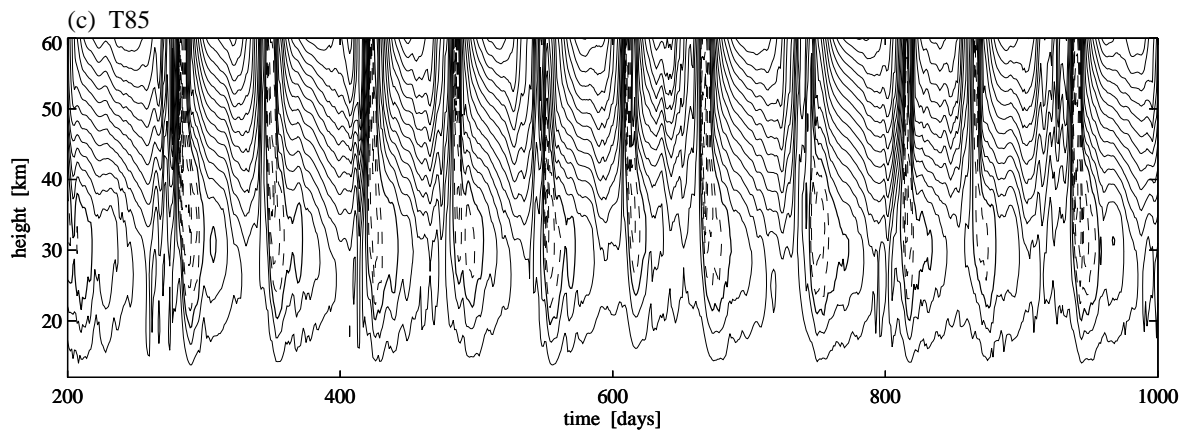


$\bar{u}(z, t)$  at 60° N

T21



T42



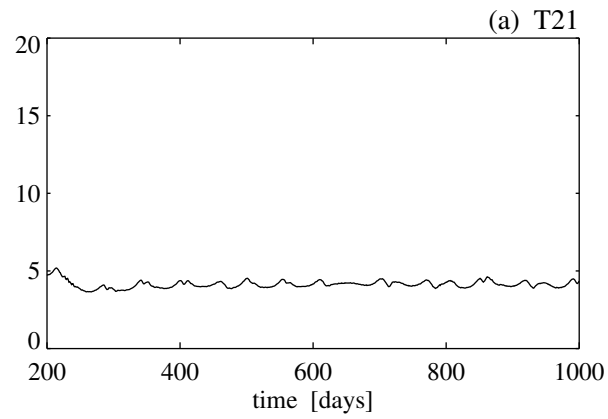
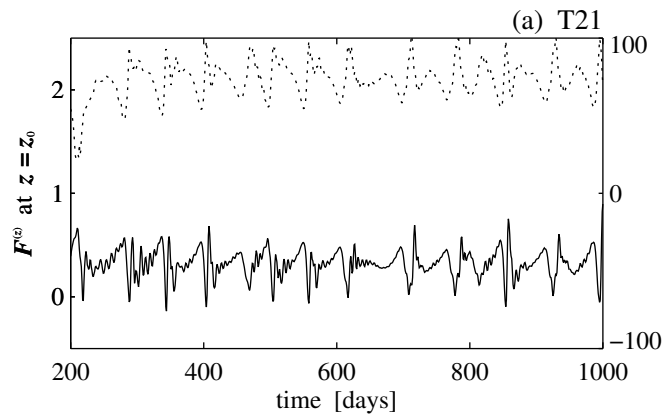
T85



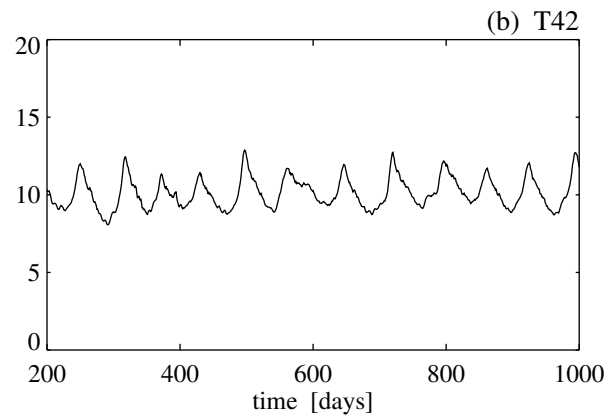
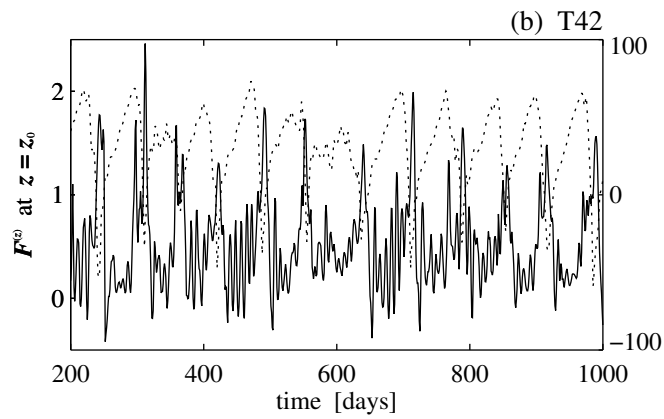
# sensitivity to horizontal resolution: T21, T42, T85

$$F(z)$$

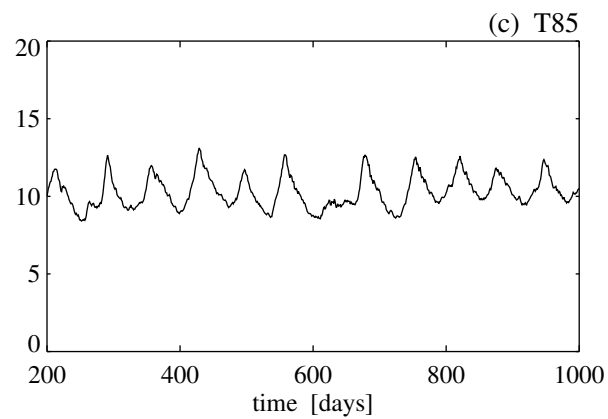
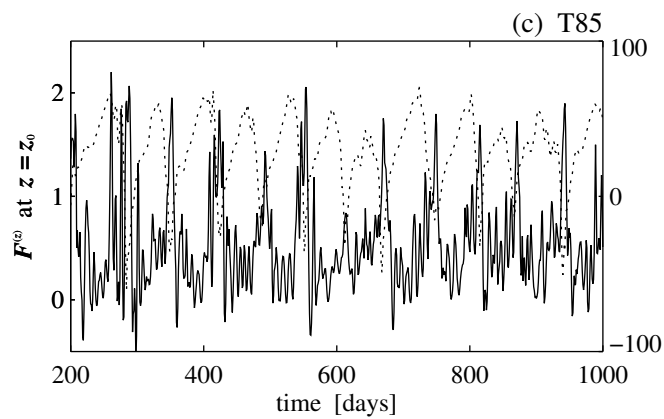
$$w^*$$



T21



T42



T85

Stratospheric response depends on forcing amplitude,  $h_0$

$h_0 \leq 200$ , steady response

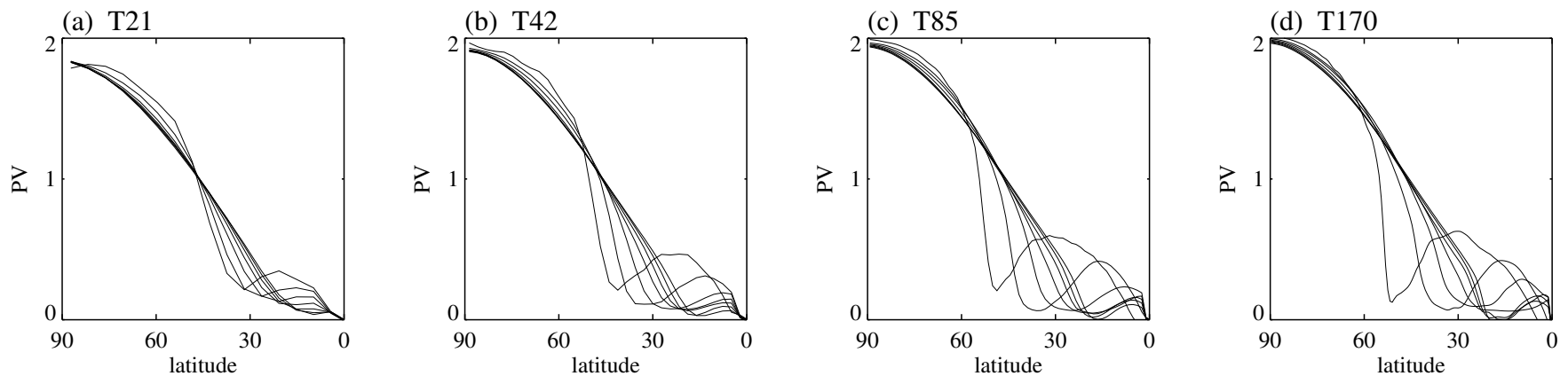
$h_0 \geq 400$ , vacillations—repeated major warmings

But critical  $h_0$  is higher at low resolution

T21: minor warmings only  $400 \leq h_0 \leq 800$

Convergence at T42 for zonal-mean quantities and time-averages;  
 $\bar{u}$ ,  $(v^*, w^*)$ ,  $\mathbf{F}$ , etc.

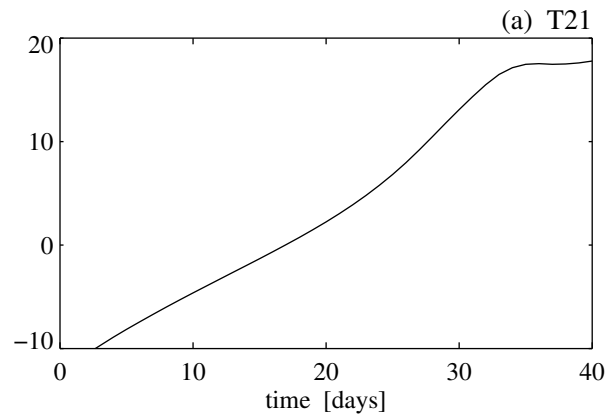
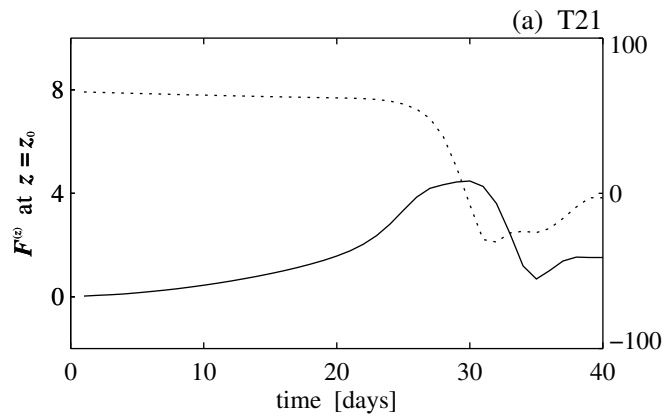
(but not instantaneous, extreme values etc.)



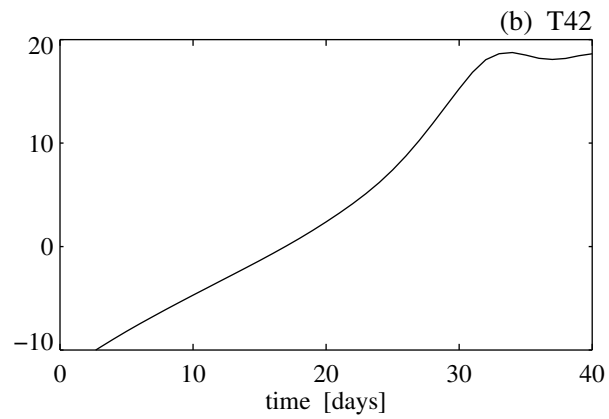
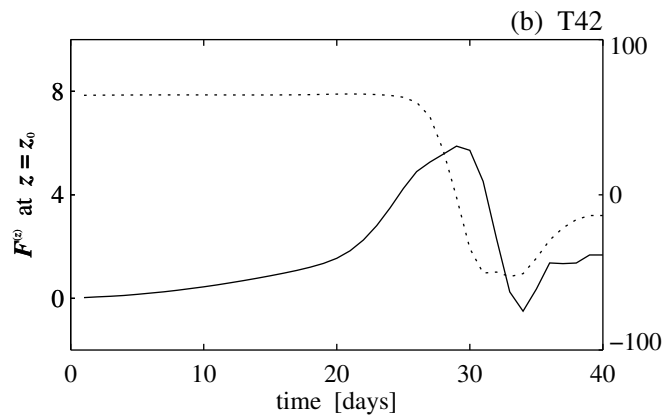
Different for initial value problem:

$$F(z)$$

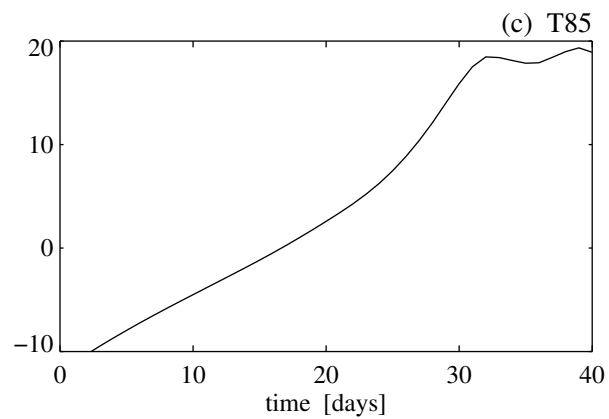
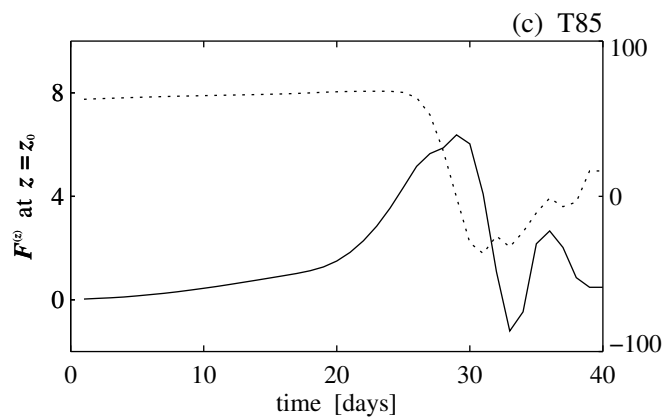
$$w^*$$



T21



T42



T85

## BAROTROPIC MODE (Esler & Scott, 2005)

Consider an idealized cylindrical polar vortex

$$Q(r, z) = \begin{cases} Q_0 & r < R \\ 0 & \text{otherwise,} \end{cases}$$

Assume disturbances to the vortex edge to position  $r = R + \eta$ :

$$\left( \frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) (\eta - r) = 0 \quad \text{on} \quad r = R + \eta$$

Get upward propagating waves (the Charney-Drazin spectrum):

$$\eta = \hat{\eta}(r) \rho_0^{-\frac{1}{2}} \exp i(k\theta + mz - \omega t) \quad \text{for } \omega \in [\omega_-, \omega_+]$$

**Also** get barotropic mode with no vertical dependence

$$\eta = \hat{\eta}(r) \exp i(k\theta - \omega t) \quad \text{for } \omega = \omega_0 < \omega_-$$

Conservation of wave activity:

$$\frac{\partial A}{\partial t} + \frac{\partial F^{(z)}}{\partial z} = 0$$

Integrate in the vertical:

$$\frac{d\mathcal{A}}{dt} = F_B^{(z)}$$

Total angular momentum:

$$\mathcal{M} + \mathcal{A} = \text{constant} \quad \implies \quad \frac{d\mathcal{M}}{dt} = -F_B^{(z)}$$

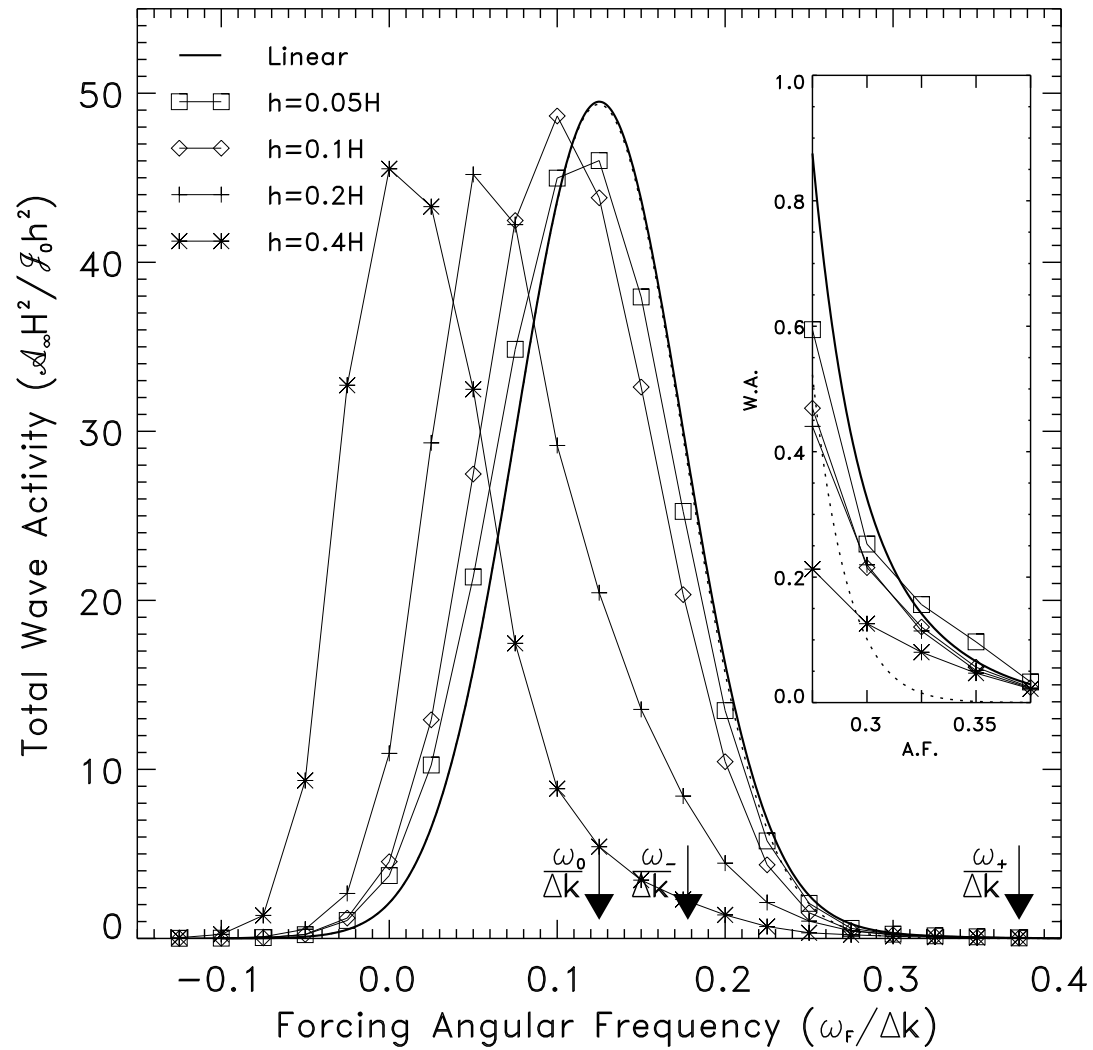
I.e.: change in vortex angular momentum = total upward wave flux

Define the relative response to the barotropic/Charney-Drazin modes:

$$\mathcal{R} = \frac{\mathcal{F}_0}{\mathcal{F}_{\text{CD}}}$$

$\mathcal{R}$  is  $\mathcal{O}(10)$ – $\mathcal{O}(100)$  ( . . . equal forcing at all frequencies)

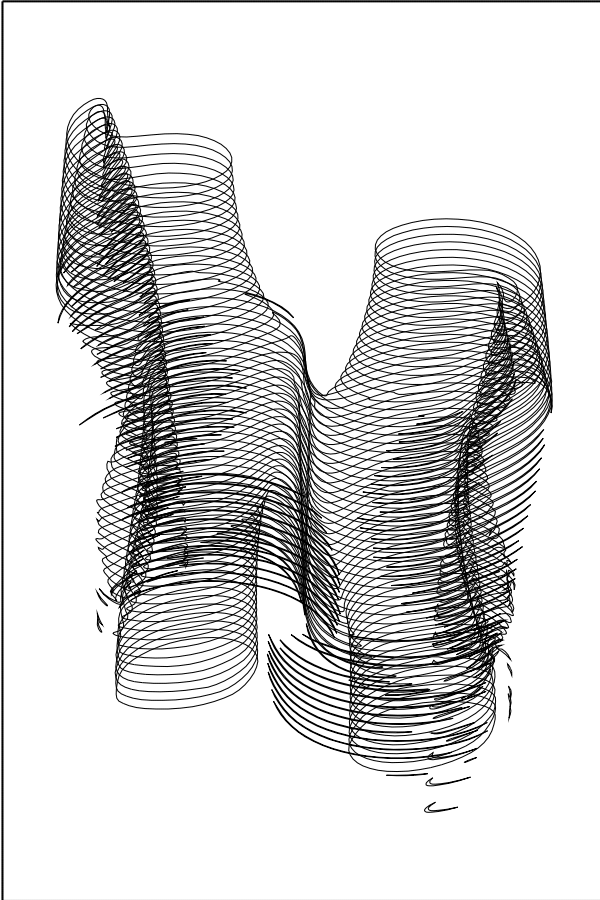
$[F_B^{(z)}]$  is given in terms of the forcing *history*



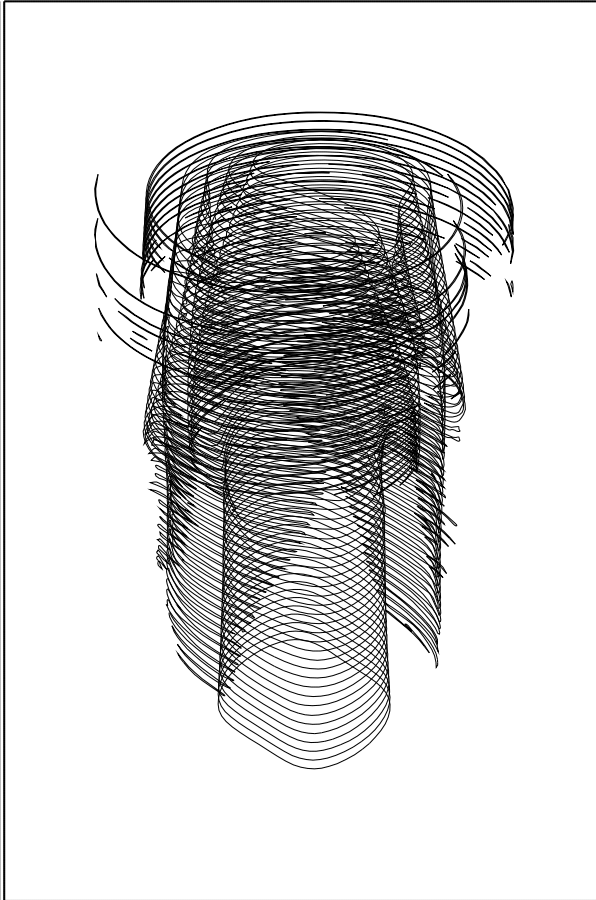
Resonant forcing at  $\omega_0$  gives rise to a much larger vortex response than forcing of upward propagating waves

Identical calculations apart from forcing frequency:

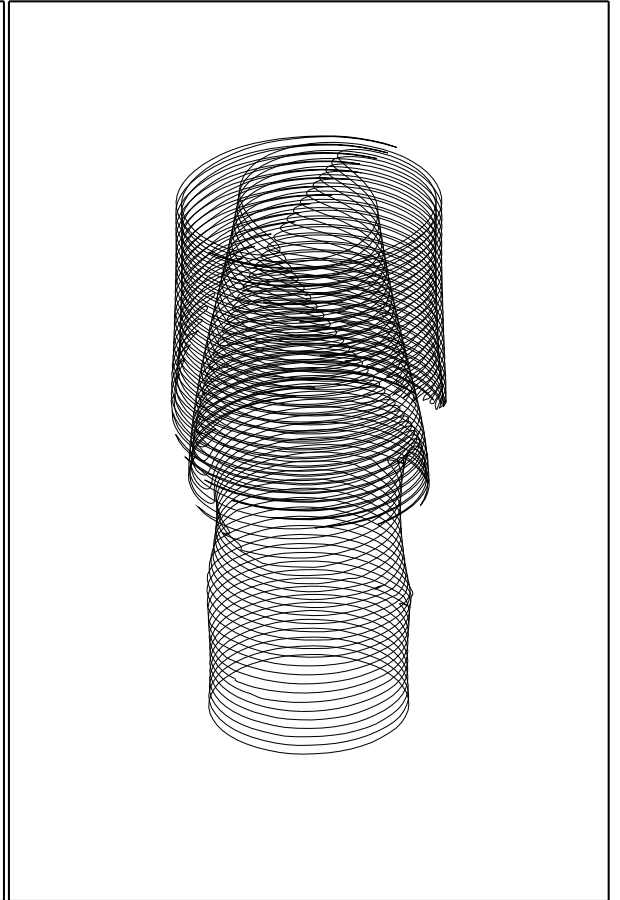
barotropic



propagating



propagating

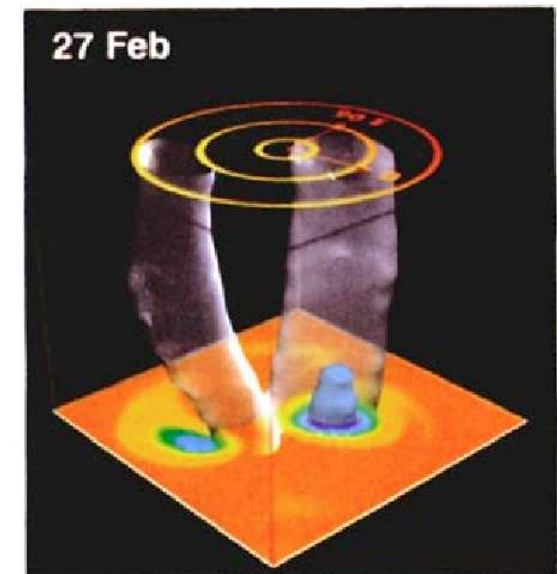
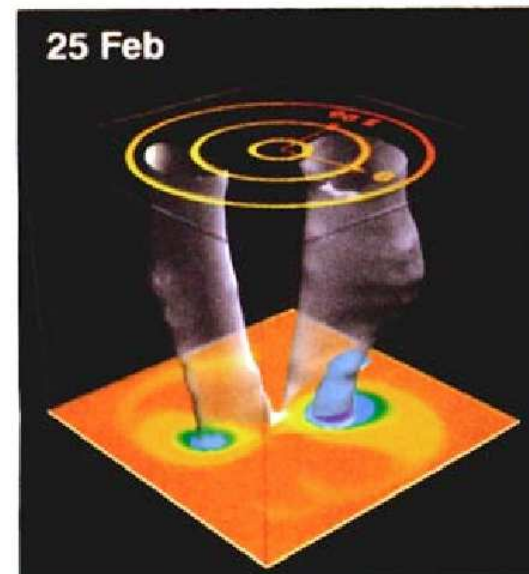
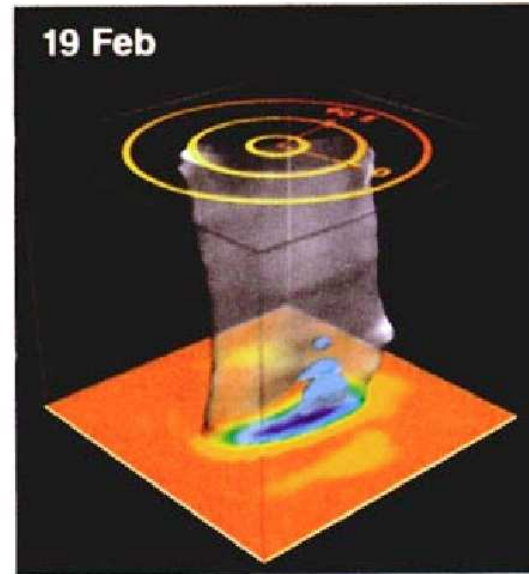
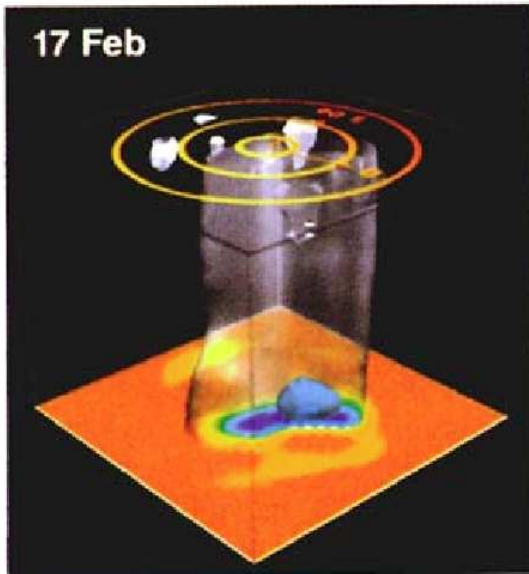


# February 1979 major warming ([Manney et al., 1994](#))

JUNE 1994

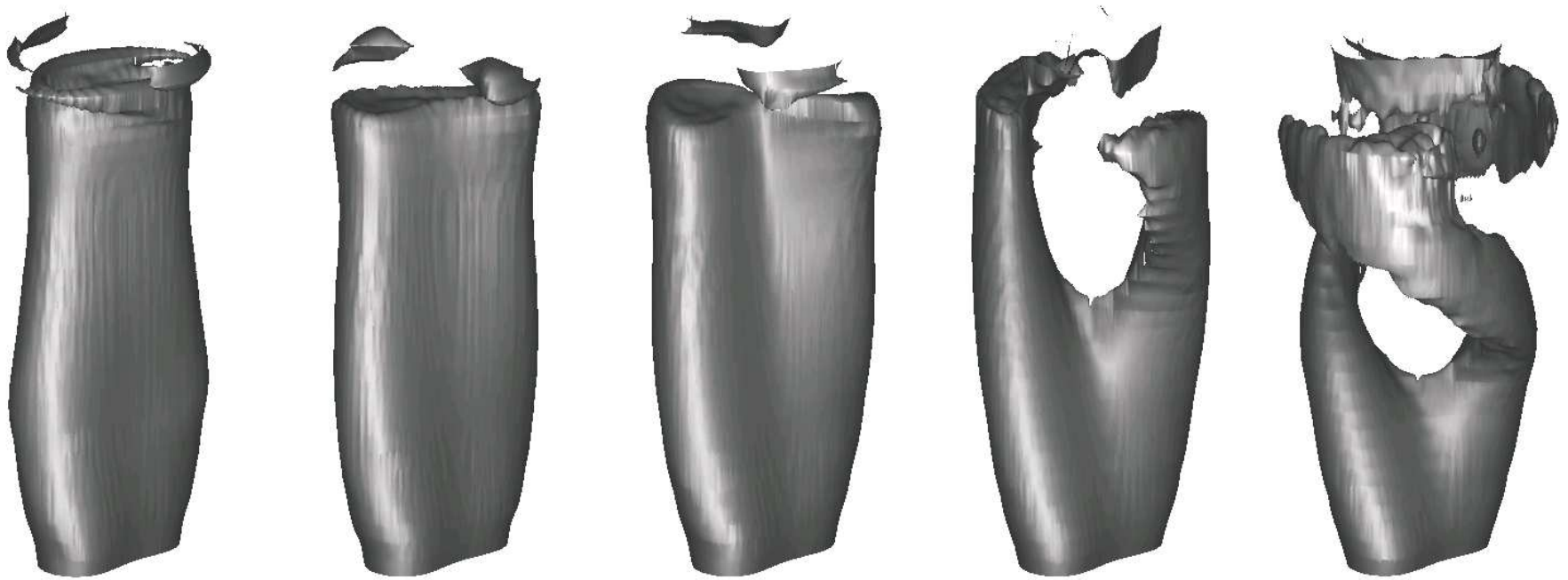
MANNEY ET AL.

1131





## Barotropic mode in internal variability



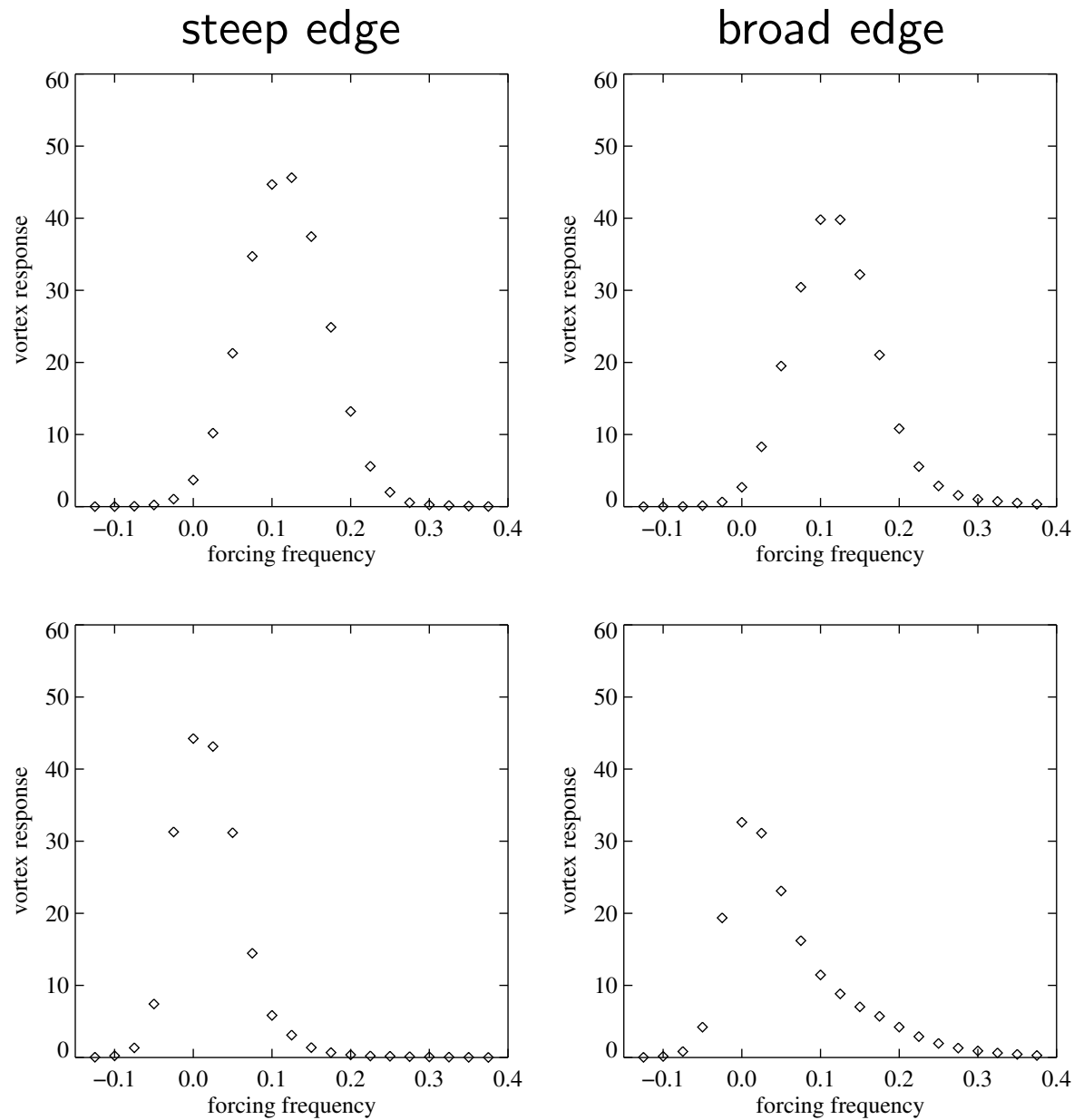
Modify traditional picture

. . . “upward propagation” of planetary waves → wave breaking, vortex deceleration, descent of critical level, etc ([Matsuno, 1971](#))

to include

. . . near-resonant excitation of the barotropic mode, in particular for vortex splitting

# Effect of resolution on excitation of the barotropic mode?



linear

nonlinear

## SUMMARY/CONCLUSIONS

For a given wave flux entering stratosphere, BD circulation is insensitive to horizontal resolution.

However, the stratosphere exerts control over upward wave fluxes.

Low resolution requires larger forcing amplitude for internal variability and repeated sudden warmings.

Or, equivalently, for a given forcing amplitude, low resolution model has fewer/weaker sudden warmings, hence weaker BD circulation:

- . . . different solution branch of internal variability
- . . . weaker response in the barotropic mode for smooth vortex edge

But T42 probably OK.